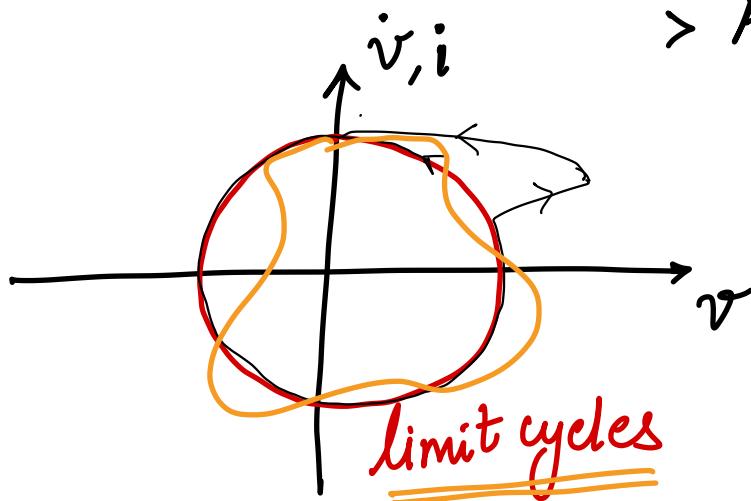




Oscillator Theory - Impulse Sensitivity Function.

(ISF model)

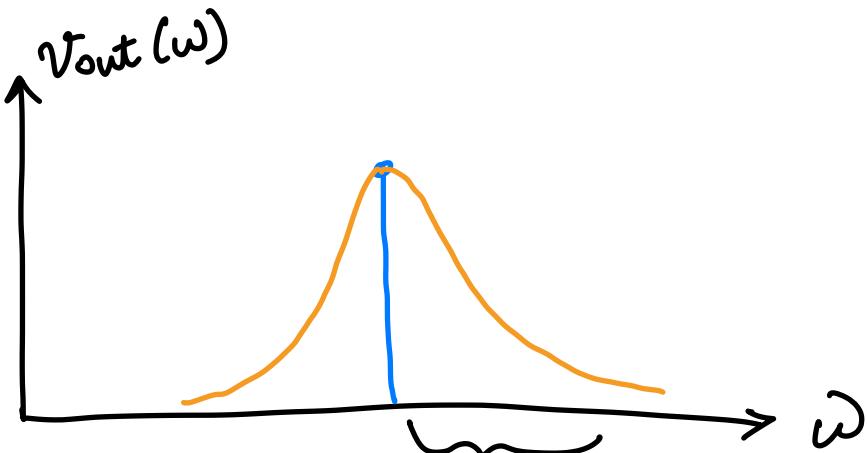


- > Amplitude fluctuations are corrected
- > Phase fluctuations persist.

$$V_{\text{out}}(t) = A \cos(\omega_0 t + \phi) \rightarrow \text{ideal}$$

$$\checkmark V_{\text{out}}(t) = A(t) \cos[\omega_0 t + \phi(t)] \rightarrow \text{real.}$$

↳ periodic with period 2π .

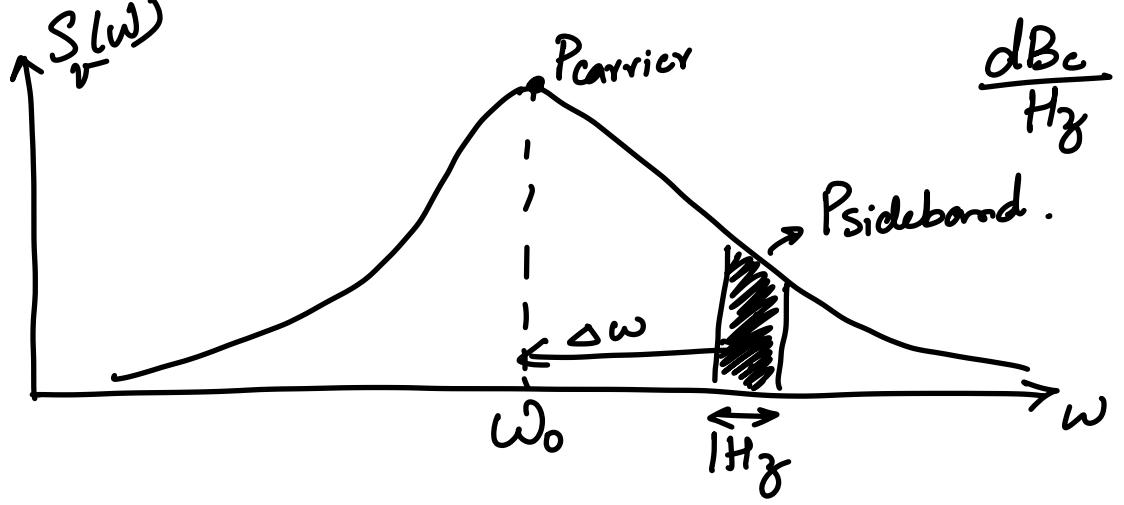


Characterized by
single sideband noise
spectral density. {Easy to measure}

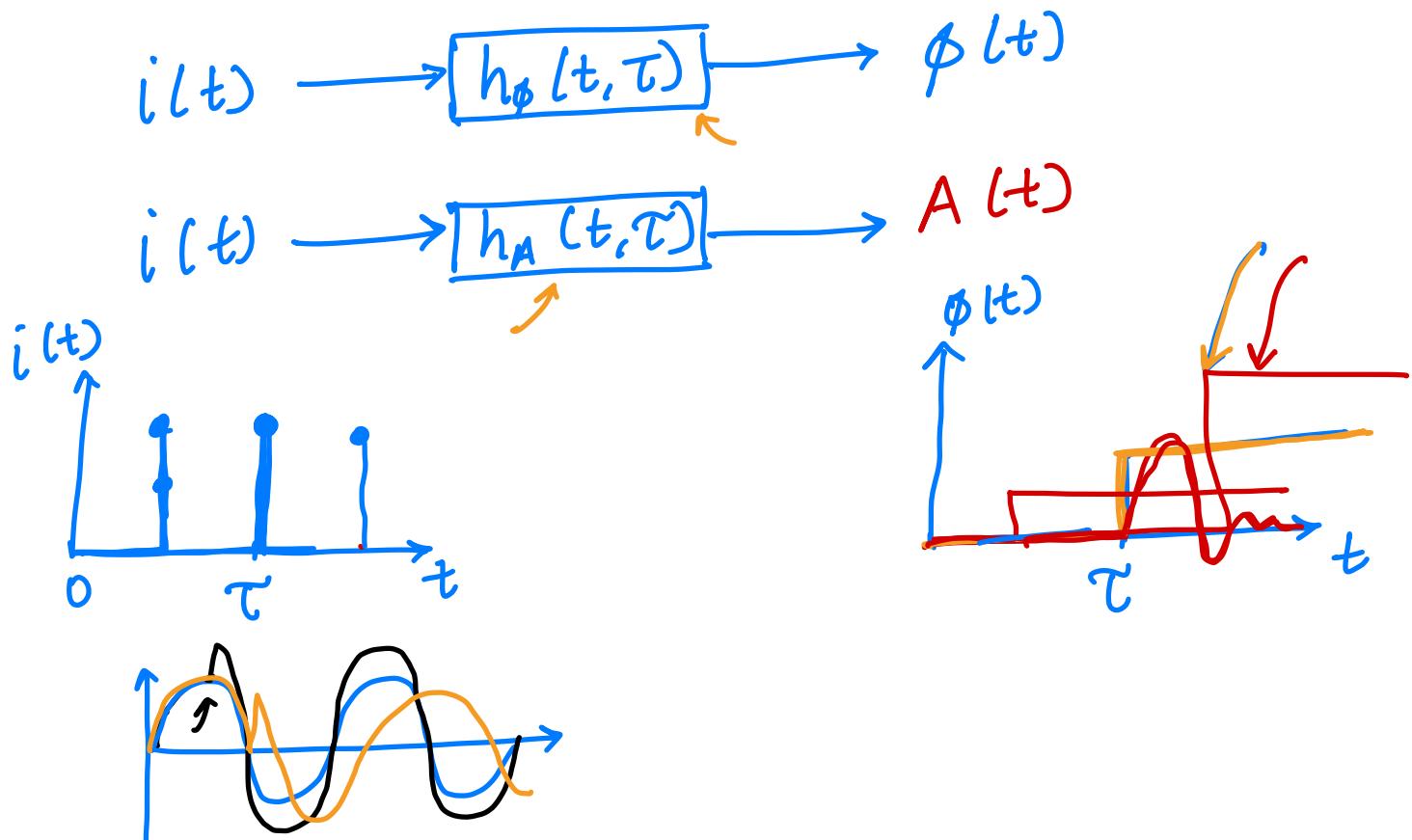
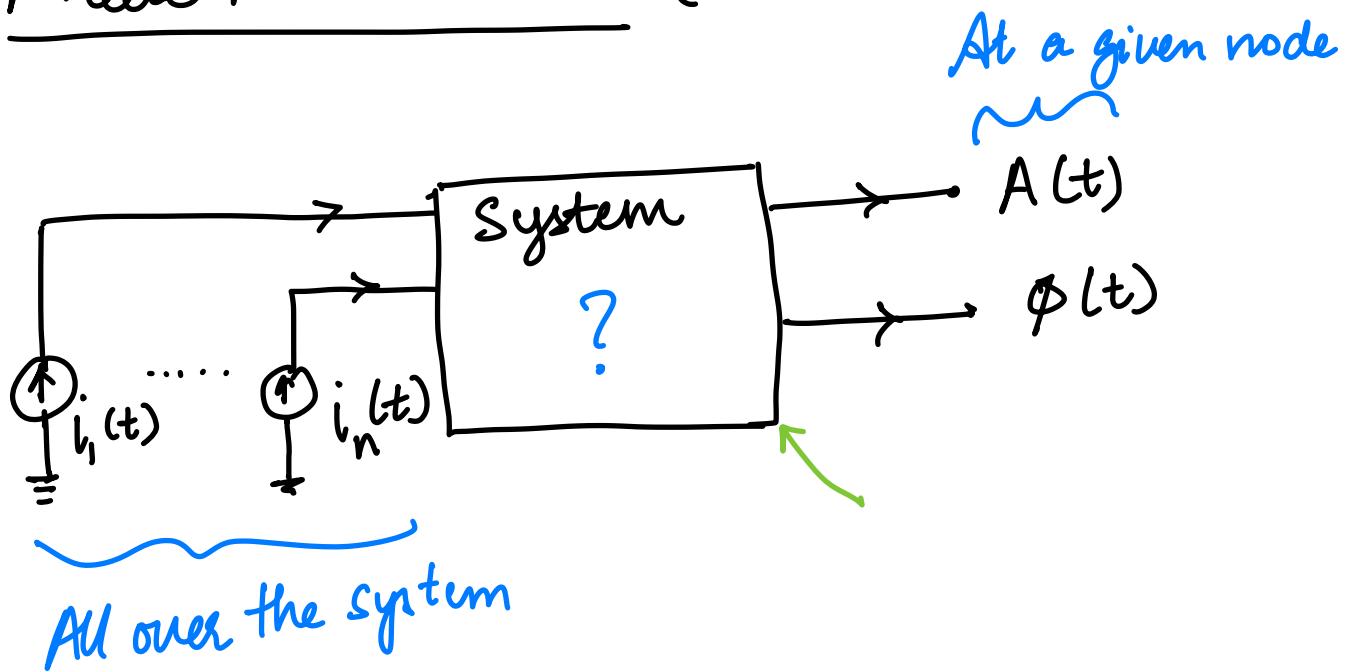
$$L_{\text{total}} \{\Delta \omega\} = 10 \log \left[\frac{P_{\text{sideband}}(\omega_0 + \Delta \omega, 1\text{Hz})}{P_{\text{carrier}}} \right]$$

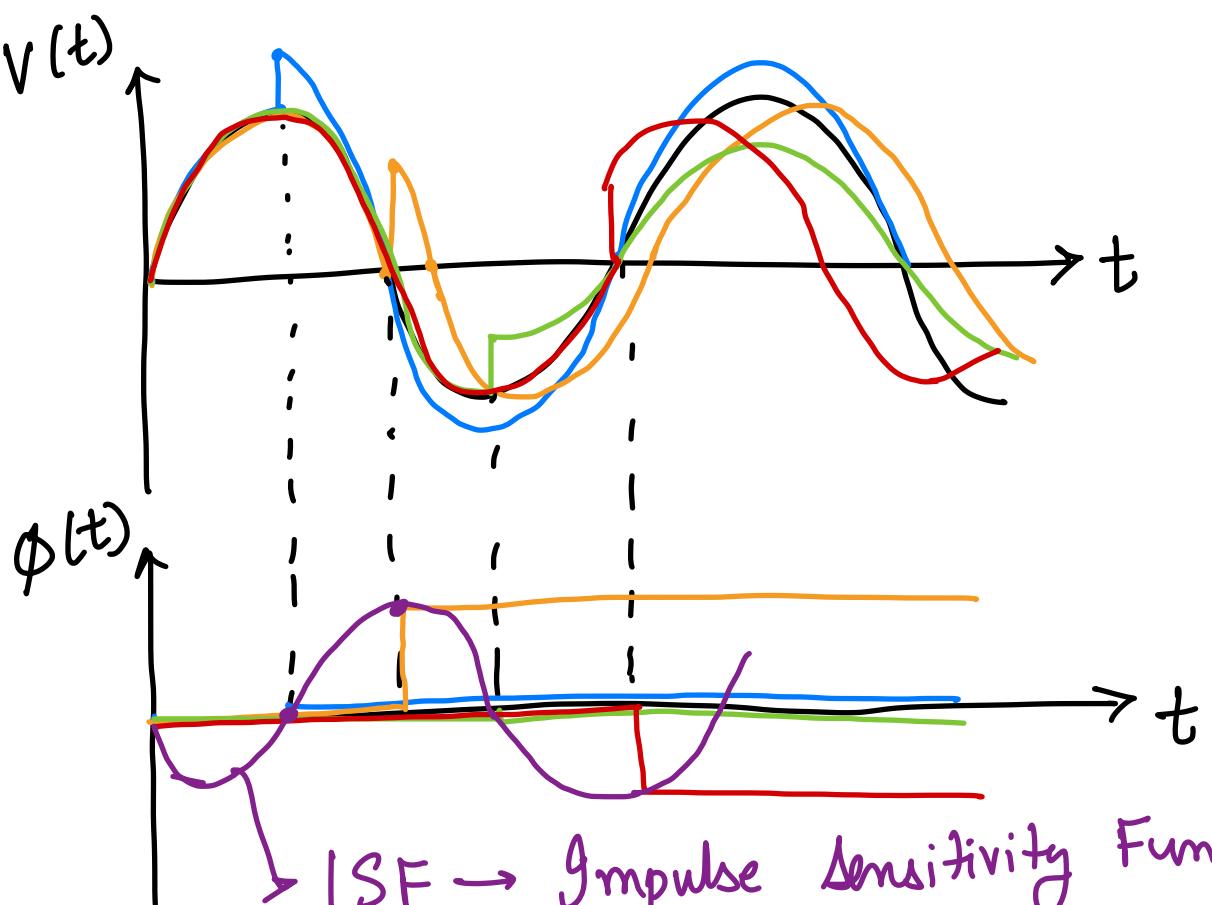
↳ Amp term
 ↳ phase term.

$\rightarrow L_{\text{phase}} \{\Delta \omega\} \rightarrow L \{\Delta \omega\} \rightarrow \underline{\text{Phase Noise.}}$



Phase Noise Model (ISF)





ISF → Impulse Sensitivity Function.

$$h_\phi(t, \tau) = \underbrace{\Gamma(\omega_0 \tau)}_{q_{\max}} u(t - \tau) \leftarrow$$

max charge displaced
on the node of interest.

Assumptions

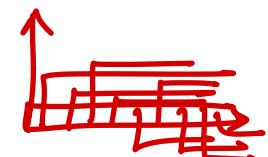
- 1) $i(t)$ $\phi(t)$ $h_\phi(t, \tau)$ is linear.
- True in practice, especially for phase noise analysis because $i(t)$ is small.
- 2) $\phi(t)$ is independent of $A(t)$ & takes the form of a step function $u(t - \tau)$.

Excess phase

$$\vec{i}(t) \xrightarrow{[h_\phi(t, \tau)]} \phi(t)$$



$$\phi(t) = \int_{-\infty}^{\infty} h_\phi(t, \tau) i(\tau) d\tau$$



$$= \frac{1}{q_{\max}} \int_{-\infty}^t \Gamma(\omega_0 \tau) i(\tau) d\tau$$

↳ Periodic

Substitute.

Fourier series expansion of Γ

$$\Gamma(\omega_0 \tau) = \frac{C_0}{2} + \sum_{n=1}^{\infty} C_n \cos(n\omega_0 \tau + \theta_n)$$

↳ limit cycle freq.

$$\Rightarrow \phi(t) = \frac{1}{q_{\max}} \left[\frac{C_0}{2} \int_{-\infty}^t i(\tau) d\tau + \sum_{n=1}^{\infty} C_n \int_{-\infty}^t i(\tau) \cos(n\omega_0 \tau) d\tau \right]$$